# A Fourier transform for all generalized functions 

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## Fourier transform in the Colombeau setting

## Definition

Let $u=\left[u_{\varepsilon}\right] \in \mathcal{G}^{s}(\Omega)$ be a CGF and $K \Subset \Omega$. Then $\int_{K} u(x) \mathrm{d} x:=\left[\int_{K} u_{\varepsilon}(x) \mathrm{d} x\right] \in \widetilde{\mathbb{R}}$. Similarly, for $\int_{\Omega} u$ if $u$ is compactly supported.

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(1) This notion of Fourier transform shares several properties with the classical one, but it lacks e.g. the Fourier inversion theorem, which holds only at the level of equality in the sense of generalized tempered distributions.
(2) We can say, the use of the multiplicative damping measure introduces a perturbation of infinitesimal frequencies that inhibit several results.

## Hyperfinite Fourier transform (HFT)

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Let $k \in^{\rho} \widetilde{\mathbb{R}}$ be a positive infinite number and let $K:=[-k, k]^{n} \subseteq{ }^{\rho} \widetilde{\mathbb{R}}^{n}$ be a functionally compact set. We define the $n$-dimensional hyperfinite Fourier transform of the GSF $f \in{ }^{\rho} \mathcal{G C}^{\infty}\left(K, \widetilde{\mathbb{C}}^{n}\right)$ as follows:

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$\left(f *_{k} g\right)(x):=\int_{\rho \widetilde{\mathbb{R}}^{n}} f(y) g(x-y) \mathrm{d} y=\int_{-k}^{k} \mathrm{~d} y_{1} \ldots \int_{-k}^{k} f(y) g(x-y) \mathrm{d} y_{n}$.

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For this type convolution we have the usual commutative, associative and distributive properties. Moreover, for $f \in{ }^{\rho} \mathcal{G C}{ }^{\infty}\left({ }^{\rho} \widetilde{\mathbb{R}},{ }^{\rho} \widetilde{\mathbb{R}}\right)$ such that $\forall x \in{ }^{\rho} \widetilde{\mathbb{R}} \exists r \in{ }^{\rho} \widetilde{\mathbb{R}}_{>0} \exists c \in{ }^{\rho} \widetilde{\mathbb{R}} \forall y \in{\overline{B^{\mathrm{E}}} r}(x) \forall j \in \mathbb{N}:\left|d^{j} f(y)\right| \leq c$ we have $f * \delta=f$, where $\delta$ is the Dirac delta distribution.

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(9) Let $k \in{ }^{\rho} \widetilde{\mathbb{R}}^{n}$. Then if $\forall x \in K$, $f\left(x_{1}, \ldots x_{j-1}, k, x_{j+1}\right)=f\left(x_{1}, \ldots x_{j-1},-k, x_{j+1}\right)=0$, we have

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We can in fact prove that inverse HFT shares the same properties as the HFT does. In the next theorem we prove that one is the inverse operation of the other.

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Let $f \in{ }^{\rho} \mathcal{G C}{ }^{\infty}\left(K,{ }^{\rho} \widetilde{\mathbb{R}}^{n}\right)$, we define $\mathcal{F}_{k}^{-1}(f)(x)=\left(\frac{1}{2 \pi}\right)^{n} \int_{K} f(\omega) e^{i x \cdot \omega} \mathrm{~d} \omega$ for all $\omega \in{ }^{\rho} \widetilde{\mathbb{R}}$. This operation is called inverse hyperfinite Fourier transform.

We can in fact prove that inverse HFT shares the same properties as the HFT does. In the next theorem we prove that one is the inverse operation of the other.

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(1) $\int_{K} \mathcal{F}_{k}(f) g(x)=\int_{K} f(x) \mathcal{F}_{k}(g)$.
(2) Hyperfinite Fourier inversion:

$$
\mathcal{F}_{k}^{-1}\left[\mathcal{F}_{k}(f)\right]=f=\mathcal{F}_{k}\left[\mathcal{F}_{k}^{-1}(f)\right]=\left(\frac{1}{2 \pi}\right)^{n} \int_{K} \mathcal{F}_{k}(f)(\omega) e^{i x \omega} \mathrm{~d} \omega
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## Examples of HFT

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(1) If $\delta(x)=\delta\left(x_{1}\right) \delta\left(x_{2}\right) \ldots \delta\left(x_{n}\right)$ is n-dimensional Dirac delta distribution with $x=\left(x_{1}, x_{2}, \ldots, x_{n}\right) \in{ }^{\rho} \widetilde{\mathbb{R}}^{n}$ and $k \geq \mathrm{d} \rho^{-a}, \forall a \in{ }^{\rho} \widetilde{\mathbb{R}}$ then $\mathcal{F}_{k}(\delta)(\omega)=1$ if $\forall \omega$ finite and $\mathcal{F}_{k}(\delta)(\omega)=0$ if $|\omega| \geq k$.

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$$
\mathcal{F}_{k}(f)(\omega)=\frac{e^{k(1-i \omega)}-e^{-k(1-i \omega)}}{1-i \omega}=\frac{\mathrm{d} \rho^{(i \omega-1)}-\mathrm{d} \rho^{(1-i \omega)}}{1-i \omega}
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## Embedding of tempered distributions


#### Abstract

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where $b \geq \mathrm{d} \rho^{-a}, a \in{ }^{\rho} \widetilde{\mathbb{R}}_{>0}$ is a linear embedding that commutes with partial derivatives and $\forall f \in \mathcal{S}\left(\mathbb{R}^{n}\right), \mathcal{F}_{k}\left(\iota_{\Omega}^{b}(f)\right)=\iota_{\Omega}^{b}(\mathcal{F}(f))$.

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$$
\mathcal{F}_{k}\left(\iota_{\Omega}^{b}(T) \cdot \mathbf{1}\right)(\omega)=\left(\iota_{\Omega}^{b}(\widehat{T})\right)(\omega), \forall \omega \in \widetilde{\Omega}_{c}, \mathbf{1}:=\mathcal{F}_{k}(\delta)
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## Examples of HFT applied in ODE

## Examples

(1) $n$-th order homogeneous generalized ODE

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\begin{aligned}
& a_{n} y^{(n)}+\ldots a_{1} y^{(1)}+a_{0} y=0, y \in{ }^{\rho} \mathcal{G} \mathcal{C}^{\infty}\left(K,{ }^{\rho} \widetilde{\mathbb{R}}\right), a_{n} \in{ }^{\rho} \widetilde{\mathbb{R}}^{*}, n \in \mathbb{N} \geq 1, \\
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(2) $n$-th order non-homogeneous generalized ODE
$a_{n} y^{(n)}+\ldots a_{1} y^{(1)}+a_{0} y=h(t), y, h \in{ }^{\rho} \mathcal{G C}^{\infty}\left(K,{ }^{\rho} \widetilde{\mathbb{R}}\right), a_{n} \in{ }^{\rho} \widetilde{\mathbb{R}}^{*}, n \in$ $\mathbb{N}_{\geq 1}, y(k)=y(-k)=0, k \in{ }^{\rho} \mathbb{R}$

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(2) $n$-th order non-homogeneous generalized ODE $a_{n} y^{(n)}+\ldots a_{1} y^{(1)}+a_{0} y=h(t), y, h \in{ }^{\rho} \mathcal{G C}^{\infty}\left(K,{ }^{\rho} \widetilde{\mathbb{R}}\right), a_{n} \in{ }^{\rho} \widetilde{\mathbb{R}}^{*}, n \in$ $\mathbb{N}_{\geq 1}, y(k)=y(-k)=0, k \in{ }^{\rho} \mathbb{R}$
(3) Generalized ODE $-u^{\prime \prime}+u=f(x), u, f \in{ }^{\rho} \mathcal{G} \mathcal{C}^{\infty}\left([-k, k],{ }^{\rho} \widetilde{\mathbb{R}}\right), u(k)=$ $u(-k)=0, k \in{ }^{\rho} \widetilde{\mathbb{R}}$.

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(9) Generalized Airy equation

$$
u^{\prime \prime}-x u=0, f \in{ }^{\rho} \mathcal{G} \mathcal{C}^{\infty}\left([-k, k],{ }^{\rho} \widetilde{\mathbb{R}}\right), u(k)=u(-k)=0, k \in{ }^{\rho} \widetilde{\mathbb{R}}
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## General procedure to apply the HFT in the study of DE

(1) We can start from a linear differential problem and assume that it has a solution $u \in{ }^{\rho} \mathcal{G C}^{\infty}\left(\widetilde{\Omega}_{c},{ }^{\rho} \widetilde{\mathbb{R}}\right)$.

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(2) We can hence take any infinite number $k \in{ }^{\rho} \widetilde{\mathbb{R}}$ and consider $K:=\{x \in \Omega| | x \mid \leq k\}, K / 2:=\{x \in \Omega| | x \mid \leq k / 2\} \subseteq{ }^{\rho} \widetilde{\mathbb{R}}^{n}$, and $\bar{u} \in{ }^{\rho} \mathcal{G C}^{\infty}\left({ }^{\rho} \widetilde{\mathbb{R}}^{n},{ }^{\rho} \widetilde{\mathbb{R}}\right)$ compactly supported in $K / 2$ and such that $\left.\bar{u}\right|_{\widetilde{\Omega}_{c}}=u$. Since $\bar{u}(x)=0$ for all $x \in\left\{x \in \widetilde{\mathbb{R}}^{n}|\forall k \in K:|x-k|>0\}\right.$, we have $\mathcal{F}_{k}\left(\partial_{j} \bar{u}\right)(\omega)=i \omega_{j} \mathcal{F}_{k}(\bar{u})(\omega)$. As usual, this allows to transform the differential problem into a simpler problem.

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(3) We finally apply the inversion theorem, at $\omega \in \widetilde{\Omega}_{c}$, so that we can recover the initial CGF $u$.

## Examples of HFT applied in PDE

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(1) Generalized wave equation
$\frac{\partial^{2} u(x, t)}{\partial t^{2}}=c^{2} \frac{\partial^{2} u(x, t)}{\partial x^{2}}, c, x \in{ }^{\rho} \widetilde{\mathbb{R}}, t \in{ }^{\rho} \widetilde{\mathbb{R}}_{\geq 0}, u \in{ }^{\rho} \mathcal{G} \mathcal{C}^{\infty}\left([-k, k] \times{ }^{\rho} \widetilde{\mathbb{R}}\right)$ with the assumptions $u(k, t)=u(-k, t)=0$ and $u_{x}(k, t)=u_{x}(-k, t)=0$ and initial conditions $u(x, 0)=f(x), u_{t}(x, 0)=g(x), x \in{ }^{\rho} \widetilde{\mathbb{R}}, f, g \in{ }^{\rho} \mathcal{G C}{ }^{\infty}\left([-k, k] \times{ }^{\rho} \widetilde{\mathbb{R}}\right)$. $u(x, t)=\frac{1}{2}[f(x+c t)+f(x-c t)]+\frac{1}{2 c} \int_{x-c t}^{x+c t} g(\xi) \mathrm{d} \xi$.

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$a^{-2} \frac{\partial u(x, t)}{\partial t}=\frac{\partial^{2} u(x, t)}{\partial x^{2}}, a, x \in{ }^{\rho} \widetilde{\mathbb{R}}, t \in{ }^{\rho} \widetilde{\mathbb{R}}_{\geq 0}, u \in{ }^{\rho} \mathcal{G C}{ }^{\infty}\left([-k, k] \times{ }^{\rho} \widetilde{\mathbb{R}}\right)$ with the analogous assumptions $u(k, t)=u(-k, t)=0$ and $u_{x}(k, t)=u_{x}(-k, t)=0$ and initial condition $u(x, 0)=f(x), x \in{ }^{\rho} \widetilde{\mathbb{R}}, f \in{ }^{\rho} \mathcal{G} \mathcal{C}^{\infty}\left([-k, k] \times{ }^{\rho} \widetilde{\mathbb{R}}\right)$ ． $u(x, t)=\frac{1}{2 \pi} \int_{-k}^{k} F(\omega) e^{-a^{2} \omega^{2} t} e^{i \omega x} \mathrm{~d} \omega, F:=\mathcal{F}_{k}(u)(\omega, 0)$ ．

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$\frac{\partial^{2} u(x, y)}{\partial x^{2}}=\frac{\partial^{2} u(x, y)}{\partial y^{2}}, x \in{ }^{\rho} \widetilde{\mathbb{R}}, y \in{ }^{\rho} \widetilde{\mathbb{R}}_{\geq 0}, u \in{ }^{\rho} \mathcal{G} \mathcal{C}^{\infty}\left([-k, k] \times{ }^{\rho} \widetilde{\mathbb{R}}\right)$ with the assumptions
$u(x, k)=0, \forall r \in{ }^{\rho} \widetilde{\mathbb{R}}_{>0}: k \geq \mathrm{d} \rho^{-r}$ and the initial condition
$u(x, 0)=f(x), x \in{ }^{\rho} \widetilde{\mathbb{R}}, f \in{ }^{\rho} \mathcal{G} C^{\infty}\left([-k, k] \times{ }^{\rho} \widetilde{\mathbb{R}}\right)$.
$u(x, y)=\frac{1}{2 \pi} \int_{-k}^{k} F(\omega) e^{-|\omega| y} e^{i \omega x} \mathrm{~d} \omega, F:=\mathcal{F}_{k}(u)(\omega, 0)$.

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(9) Paley-Wiener theorem using hyperfinite series.

## Contact

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Thank you for your attention!

